

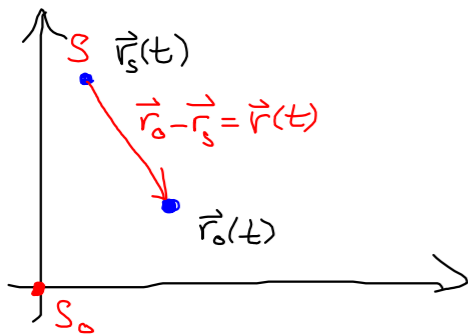
# Lecture 07.1 - Physics 523

Announcements Midterm exam next tuesday, in Lecture! BE THERE!

## Non-inertial Frames

Recall last week we discussed physics in 2 different frames:  $S_0$  (inertial) +  $S$  (non-inertial)

Classical mech, in frame  $S_0$  obeys Newton's Law  $\vec{F} = m\ddot{\vec{r}}_0$ . If frame  $S_0$  sees trajectory  $\vec{r}_0(t)$ , & frame  $S$  moves through trajectory  $\vec{r}_S(t)$ , then  $S$  sees trajectory  $\vec{r}_0 - \vec{r}_S \equiv \vec{r}(t)$



$$\vec{F} = m\ddot{\vec{r}}_0 = m(\ddot{\vec{r}} + \ddot{\vec{r}}_S)$$

$$m\ddot{\vec{r}} = \vec{F} - m\ddot{\vec{r}}_S$$

$$\boxed{m\ddot{\vec{r}} = \vec{F} - m\vec{A}}$$

$\uparrow$  acceleration of N-I frame  $S$

We can now utilize this formula for addressing physics as observed in accelerating reference frames!

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# Tidal Forces

The earth's oceans respond to the relative orientations of the earth + moon

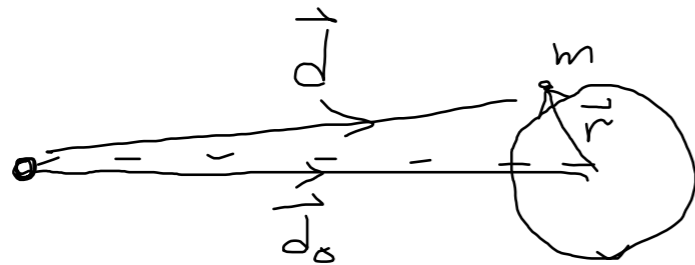


Forces on mass  $m$  on surface of earth

1)  $-mg\hat{r}$  earth's center is origin

2)  $F_{\text{moon}} = -\frac{GM_0 m}{d^2} \hat{d}$

3) Other local forces,  $\vec{F}_{ng}$



Earth has centripetal accel,  $\vec{A} = -\frac{GM_0}{d_0^2} \hat{d}_0$

Using non-inertial frame mod. of Newton's 2nd  $m\ddot{\vec{r}} = \vec{F} - m\vec{A}$ :  $m\ddot{\vec{r}} = -mg\hat{r} - GM_0 m \frac{\hat{d}}{d^2} + \vec{F}_{ng} + \frac{GM_0 m}{d_0^2} \hat{d}_0$

$m\ddot{\vec{r}} = -mg\hat{r} + \vec{F}_{ng} + \vec{F}_{\text{tidal}}$

$$\vec{F}_{\text{tidal}} = -GM_0 m \left( \frac{\hat{d}}{d^2} - \frac{\hat{d}_0}{d_0^2} \right)$$

## Equipotential Surface

Claim: The oceans surface is an equipotential surface

Proof for conservative forces  $\vec{F} = -\nabla U$  for  $\hat{n}$  on surface,  $\vec{F} \cdot \hat{n} = -(\nabla U) \cdot \hat{n}$  if surface is in equilibrium,  $\vec{F} \cdot \hat{n} = 0$   
(no lateral motion of water)

So  $\nabla U$  vanishes along surface  $\Rightarrow$  equipotential

$$\vec{F}_{\text{tidal}} = -GM_0 m \left( \frac{\hat{d}}{d^2} - \frac{\hat{d}_0}{d_0^2} \right) \quad \text{like the gradient of 2 terms} \quad 1) \sim -\frac{1}{d} \quad 2) -\frac{x}{d_0^2}$$

$$U_{\text{tidal}} = -GM_0 m \left( \frac{1}{d} + \frac{x}{d_0^2} \right) \\ = -\frac{GM_0 m}{d_0} \left( \frac{d_0}{d} + \frac{x}{d_0} \right)$$

$$U_{\text{tot}} = mgh - GM_0 m \left( \frac{1}{d} + \frac{x}{d_0^2} \right) \stackrel{\text{low tide}}{=} -\frac{GM_0 m}{d_0} \left( \frac{d_0}{d} \right) = -\frac{GM_0 m}{d_0} \left( \frac{d_0}{\sqrt{R_e^2 + d_0^2}} \right) \stackrel{\text{Taylor exp}}{=} -\frac{GM_0 m}{d_0} \left( 1 - \frac{1}{2} \left( \frac{R_e}{d_0} \right)^2 \right) \\ \stackrel{\text{height above low tide}}{=} mgh - \frac{GM_0 m}{d_0} \left( \frac{d_0}{d} + \frac{x}{d_0} \right) = mgh - \frac{GM_0 m}{d_0} \left[ \frac{d_0}{d_0 \pm R_e} \pm \frac{R_e}{d_0} \right] \\ \stackrel{\text{1st order term cancels}}{\approx} mgh - \frac{GM_0 m}{d_0} \left[ 1 + \left( \frac{R_e}{d_0} \right)^2 \right]$$

$$mgh = \frac{GM_0 m}{d_0} \cdot \frac{3}{2} \frac{R_e^2}{d_0^2}$$

$$\text{and } g = \frac{GM_e}{R_e^2}$$

$$H = \frac{3}{2} \frac{M_0}{M_e} \frac{R_e^4}{d_0^3}$$

All this came from  $m\ddot{\vec{r}} = \vec{F} - m\vec{A}$

Example

Spinning w/ a coffee cup  
angular frequency  $\omega$

$$m\ddot{\vec{r}} = \vec{F} - m\vec{A}$$

$$\vec{F} = m\vec{g} + \vec{F}_{ng} \quad ; \quad \vec{A} = -v^2 \frac{\hat{r}}{r}$$

$$m\ddot{\vec{r}} = m\vec{g} + \frac{mv^2}{r}\hat{r} + \vec{F}_{ng} = -\vec{\nabla}U_{eff} + \vec{F}_{ng}$$

$$-\vec{\nabla}U_{eff} = m\vec{g} + \frac{mv^2}{r}\hat{r} = m\vec{g} + m\omega^2 r \hat{r}$$

$U_{eff} = mgh - m\omega^2 r^2 = \text{const} = 0$  (surface of liquid is equipotential)  
( $h=0$  for  $r=0$  convention)

$$h = \frac{\omega^2 r^2}{g}$$