

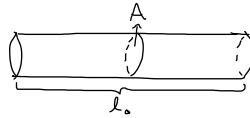
# Lecture 3.1 "Oscillations are everywhere" off-diagonal terms in $U(x,y)$ B-meson oscillation

Two crucial features: 1) Inertial component  $\rightarrow$  Carries kinetic energy "Mass" 2) Elastic component  $\rightarrow$  Stores elastic potential energy "Spring" } These are much more general than just

Example Mass  $m$  on finite mass spring w/ mass per length  $\lambda$ , length  $L$  (actually quite complicated to analyze in detail!)

## Example Vibrations of semi-rigid bodies

Consider rod w/ length  $l_0$ , cross-sectional area  $A$ :  
(or wire)



Consider stretching rod/wire by  $\Delta l \ll l_0$  (Applying force  $\Delta P$  to end)

"Strain"  $\equiv \Delta l / l_0$

"Stress"  $\equiv \frac{\Delta P}{A}$  units of pressure

For small displacements (strain  $= \frac{\Delta l}{l_0} \ll 1$ )

$\frac{\text{stress}}{\text{strain}} = \text{constant}$

↑  
Young's Modulus of Elasticity

Property of Material!

Equation for stretching of material

$$\left(\frac{dF/A}{d(\Delta l/l_0)}\right) = -Y, \text{ or } dF = -\frac{AY}{l_0} d\Delta l$$

$\Rightarrow$  for uniform stretching by amount  $x$ :  $F = -\frac{AY}{l_0} x$  like Hooke's Law!

Note: This fails @ strains larger than  $\sim 0.1\%$  1) non-linear terms

2) wire breaks

Analyzing oscillations:

Hang mass  $m$  on wire  $\rightarrow F = -\frac{AY}{l_0} x = m\ddot{x} \Rightarrow \ddot{x} = -\frac{AY}{l_0 m} x \Rightarrow \omega^2 = \frac{AY}{l_0 m}$

$\omega = 2\pi f = \frac{2\pi}{T} \Rightarrow T = 2\pi \sqrt{\frac{l_0 m}{AY}}$

for  $l_0 = 1\text{m}, m = 1\text{kg}$

steel w/ diameter 1mm  $T = 16 \times 10^{-2} \text{sec}$   
 $f = 60\text{Hz}$

## DEMO-TORSION PENDULUM - Also book examples

1) Floating objects  $m\ddot{y} = -\rho g A y$

$\rightarrow$  buoyant force Archimedes!

2) Complicated pendula  $I\ddot{\theta} = -mgh\theta$

$\uparrow$  dist. to CM

3) Water in tube  $\rho A l \ddot{y} = 2\rho g A y$

Trick: I.D. the KE and PE of system conserve  $E_{\text{tot}}$ , take derivative!

In-class exercise: Show that  $Ae^{pt}$  solves  $m\ddot{x} = -kx$  if  $p = i\omega = i\sqrt{\frac{k}{m}}$

Hint  $\frac{d}{dx} e^{\lambda x} = \lambda e^{\lambda x}$

Now, what is  $p$  for  $m\ddot{x} + b\dot{x} + kx = 0$ ?

Damping In nearly all cases, oscillating system eventually lose energy to the environment (heat, other waves, etc.)

Motion of body through fluid  $\rightarrow$  Resistive force  $R(v) = \underbrace{bv}_{\rightarrow \text{first non-trivial term}} + cv^2 + \dots$  for small  $v$

Newton's Law  $F = -kx - \underbrace{b\dot{x}}_{bv} = m\ddot{x}$  Exactly the math example you just worked out!  $x \rightarrow z(t) = Ae^{pt}$

$$mp^2x + bp\dot{x} + kx = 0$$

$$p = \frac{-b \pm \sqrt{b^2 - 4km}}{2m} = -\frac{b}{2m} \pm i \sqrt{\frac{k}{m} - \left(\frac{b}{2m}\right)^2}$$

Solution  $z(t) = A_+ e^{(-\frac{b}{2m} + i\omega)t} + A_- e^{(-\frac{b}{2m} - i\omega)t}$  Oscillations w/ exponentially decaying amplitude!

Equiv.  $x(t) = \text{Re}[A e^{-\frac{b}{2m}t} e^{i(\omega t + \alpha)}]$